# Blow-up in non-autonomous semilinear pseudoparabolic equations

#### Sujin Khomrutai, Nataphan Kitisin\*

Department of Mathematics and Computer Science, Faculty of Science, Chulalongkorn University, Bangkok 10330 Thailand

\*Corresponding author, e-mail: nataphan.k@chula.ac.th

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**ABSTRACT**: We study the blow-up property for weak solutions to the Cauchy problem of non-autonomous semilinear pseudoparabolic equations. Given the growth bound of the non-autonomous coefficient, the Fujita-type critical exponent is obtained.

KEYWORDS: critical exponent, test function method, global solutions

## INTRODUCTION

We study non-negative weak solutions u = u(x, t) of the Cauchy problem

$$\begin{cases} \partial_t u - \Delta \partial_t u = \Delta u + V(x)u^p, & \text{in } \mathbb{R}^n \times \mathbb{R}_+, \\ u(x,0) = u_0(x), & \text{in } \mathbb{R}^n \end{cases}$$
(1)

where p > 1 is a constant and  $V, u_0$  are given nonnegative functions. This partial differential equation (PDE) is called a *pseudoparabolic equation*<sup>1-4</sup>. It is semilinear and non-autonomous owing to the coefficient V(x) on the right-hand side. Nonlinear pseudoparabolic equations have been proposed to model many physical systems; for instance, the non-steady flow of second-order fluids in one space dimension<sup>5</sup>, seepage of homogeneous fluids through fissured rock<sup>6</sup>, heat conduction involving two temperatures<sup>7</sup>.

The equation (1) is also closely related with the following non-autonomous semilinear heat equation

$$\partial_t u = \Delta u + V(x)u^p$$

and the latter has been widely investigated by many authors<sup>8</sup>. In the case  $V(x) \equiv 1$ , the problem (1) becomes autonomous and was investigated by Cao et al<sup>9</sup>. In their paper, the existence of mild solutions, which are also weak solutions, was established. Using the energy method, the authors obtained the critical exponent of the problem, denoted by  $p_c$ , for the class of classical solutions:

$$p_{\rm c} = 1 + \frac{2}{n}.\tag{2}$$

This means that if 1 then every nontrivialnon-negative solution to the problem blows up in some $finite time <math>T_0$ , i.e.,  $\lim_{t\to T_0^-} ||u(\cdot,t)||_{L^{\infty}} = \infty$ . On the other hand, if  $p > p_c$  there are both blowing-up solutions (for sufficiently large  $u_0$ ) and global-in-time solutions (for sufficiently small  $u_0$ ). Even though the blow-up phenomenon has played an important role in PDE theory<sup>10,11</sup>, the blowing-up problem of (1) for non-constant V, however, remains open.

In this study, the Cauchy problem (1) with a broader class of functions V is considered and the critical exponent analogous to (2) is obtained. The energy method does not seem to work for the weak solution in the case where V is non-constant and therefore a new approach is needed. The technique employed here is the test function (or nonlinear capacity) method<sup>12</sup>. Other important related questions (e.g., existence, uniqueness, regularity, and large-time asymptotic) will be addressed in our forthcoming papers.

### PRELIMINARIES

Let 
$$Q_T = \mathbb{R}^n \times [0, T)$$
 and  $Q_\infty = \mathbb{R}^n \times [0, \infty)$ .

**Lemma 1 (Folland <sup>13</sup>)** Let  $(X, \mu)$  be a measure space and  $1 \le p, q \le \infty$ .

(i) For all  $a, b \ge 0$  and  $\lambda \in [0, 1]$ , we have

$$a^{\lambda}b^{1-\lambda} \leqslant \lambda a + (1-\lambda)b.$$

(ii) If  $f \in L^p(X)$ ,  $g \in L^q(X)$  where 1/p + 1/q = 1, then  $h = fg \in L^1(X)$  and

$$||h||_{L^1(X)} \leq ||f||_{L^p(X)} ||g||_{L^q(X)}.$$

The solutions of (1) considered in this paper are weak solutions which are defined as follows.

**Definition 1** A function u is called a *weak* (or *distributional*) solution to the problem (1) on I = [0, T) provided

(i)  $u \in C(I; L^1_{loc}(\mathbb{R}^n)), Vu^p \in L^1_{loc}(I; L^1_{loc}(\mathbb{R}^n)),$ (ii) for all  $\varphi \in C^3_c(Q_T)$ , the following identity holds:

$$\iint_{Q_T} u(\partial_t \varphi - \Delta \partial_t \varphi + \Delta \varphi) + \iint_{Q_T} V u^p \varphi$$
$$= \int_{\mathbb{R}^n} u_0(\Delta \varphi - \varphi)|_{t=0} \quad (3)$$

If (i) and (ii) are true with  $T = \infty$ , then u is called a global weak solution.

The following lemma<sup>12</sup> will be used to construct the test functions needed below. For the completeness of the paper, the proof is provided.

**Lemma 2** For any  $q \in (1, \infty)$ , there is a  $C^3$  function  $\phi : \mathbb{R} \to [0, 1]$  with  $\phi(s) = 1$  if  $s \leq 1, 0 \leq \phi(s) \leq 1$  if  $1 \leq s \leq 2, \phi(s) = 0$  if  $s \geq 2$ , and

$$|\phi'(s)|^{q} + |\phi''(s)|^{q} + |\phi'''(s)|^{q} \leq C_{\phi}\phi(s)^{q-1}$$

for all  $s \in \mathbb{R}$ , for some constant  $C_{\phi} > 0$ .

*Proof*: Choose a function  $\zeta \in C^3(\mathbb{R}, [0, 1])$  with

$$\zeta(s) \begin{cases} = 1, & s \leq 1, \\ \in (0,1), & 1 < s < 2, \\ = 0, & s \ge 2 \end{cases}$$

and let  $\phi(s) = \zeta(s)^{3q}$ . Then

$$\begin{aligned} |\phi'|^{q} &= |3q\zeta'|^{q} \zeta^{(3q-1)q} \leqslant C\phi^{q-1} \\ |\phi''|^{q} &= |3q(3q-1)(\zeta')^{2} + 3q\zeta\zeta''|^{q} \zeta^{(3q-2)q} \\ &\leqslant C\phi^{q-1} \\ |\phi'''|^{q} &= |3q(3q-1)(3q-2)(\zeta')^{3} \\ &+ 3(3q)(3q-1)\zeta\zeta'\zeta'' + 3q\zeta^{2}\zeta'''|^{q}\zeta^{(3q-3)q} \\ &\leqslant C\phi^{q-1} \end{aligned}$$

because  $(3q - i)q \ge 3q(q - 1)$  for i = 1, 2, 3 and  $0 \le \zeta \le 1$  where C > 0 is a constant depending only on  $q, \|\zeta'\|_{L^{\infty}}, \|\zeta''\|_{L^{\infty}}$ , and  $\|\zeta'''\|_{L^{\infty}}$ .  $\Box$ 

**Remark 1** Generally, for any  $q \in (1, \infty)$  and  $k \in \mathbb{N}$ , there is  $\phi \in C^k([0, \infty), [0, 1])$  satisfying

$$\sum_{i=1}^{k} \left| \phi^{(i)}(s) \right|^{q} \leqslant C \phi(s)^{q-1} \quad \forall s \ge 0$$

for some constant C.

### MAIN RESULTS

We will consider the class of functions V in (1) that satisfies the following assumption.

Assumption. The function V(x) has an order of growth of at least  $\sigma > -2$ , in the sense that there exists  $x_0 \in \mathbb{R}^n$  and a constant  $c_0 > 0$  such that

$$V(x) \ge c_0 \left| x - x_0 \right|^{\sigma} \tag{4}$$

for almost every  $x \in \mathbb{R}^n$ .

The proof below is valid for arbitrary  $x_0$ . However, for simplicity of presentation and without loss of generality, we let  $x_0 = 0$ .

**Theorem 1** Assume (4) on V and let  $1 + (\sigma^+)/n where <math>\sigma^+ = \max\{0, \sigma\}$ . If  $0 \leq u_0 \in L^1(\mathbb{R}^n)$  with  $||u_0||_{L^1(\mathbb{R}^n)} > 0$ , then there is no nontrivial, non-negative global weak solution u to the problem (1).

*Proof*: The theorem will be proved by contradiction. We therefore assume the contrary that the problem (1) admits a non-trivial, global weak solution u. We divide the proof into 5 steps.

Step 1. Define the operator

$$\mathcal{A}\varphi := \partial_t \varphi - \Delta \partial_t \varphi + \Delta \varphi,$$

for all test functions  $\varphi \in C_c^3(Q_\infty)$ . Then u satisfies for all  $\varphi$  the identity

$$\iint_{\operatorname{supp}\mathcal{A}\varphi} u\mathcal{A}\varphi + \iint_{\operatorname{supp}\varphi} Vu^p \varphi = \int_{\mathbb{R}^n} u_0(\Delta\varphi - \varphi)|_{t=0}.$$
(5)

Choose  $\varphi$  to satisfy  $0 \leq \varphi \leq 1$  and  $\varphi|_{\text{supp } \mathcal{A}\varphi} > 0$  a.e. so that  $\text{supp } \mathcal{A}\varphi \subset \text{supp } \varphi$ .

Let  $K = \operatorname{supp} \varphi$  and  $K' = \operatorname{supp} \mathcal{A} \varphi$ . By the Hölder and Young inequalities, we have

$$\iint_{K'} |u\mathcal{A}\varphi| \\
\leqslant \left(\iint_{K'} Vu^{p}\varphi\right)^{1/p} \left(\iint_{K'} \frac{|\mathcal{A}\varphi|^{q}}{(V\varphi)^{q-1}}\right)^{1/q} \\
\leqslant \frac{1}{p} \iint_{K} Vu^{p}\varphi + \frac{1}{q} \iint_{K'} \frac{|\mathcal{A}\varphi|^{q}}{(V\varphi)^{q-1}}, \quad (6)$$

where  $q = (p/(p-1)) \in (1, \infty)$ . Combining (5), (6) with the assumption  $V(x) \ge c_0 |x|^{\sigma}$  a.e. yields the estimate

$$\iint_{K} V u^{p} \varphi \leqslant \frac{1}{c_{0}^{q-1}} \iint_{K'} \frac{|\mathcal{A}\varphi|^{q}}{(|x|^{\sigma} \varphi)^{q-1}} + q \int_{\mathbb{R}^{n}} u_{0}(\Delta \varphi - \varphi)|_{t=0}.$$
 (7)

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**Step 2.** Let us further specify  $\varphi$ . Fix a function  $\phi$  satisfying Lemma 2. For  $R \gg 1$ , define  $\varphi$  by

$$\varphi(x,t) = \phi\left(\frac{t+|x|^2}{R^2}\right).$$

Below, the following rescaling variables will be used:  $\tau = t/R^2$ ,  $\xi = |x|/R$ , and  $s = \tau + \xi^2$ . As subsets in the  $\xi$ t-plane,  $\{(\xi, \tau) : 1 \leq \tau + \xi^2 \leq 2\} = K' \subset K$  and  $K \subset [0, \sqrt{2}] \times [0, 2]$ . Direct computation shows that

$$\begin{split} \partial_t \varphi &= \frac{1}{R^2} \phi', \\ \Delta \varphi &= \frac{1}{R^2} \Big( 4\xi^2 \phi'' + 2n\phi' \Big), \\ \partial_t \Delta \varphi &= \frac{1}{R^4} (4\xi^2 \phi''' + 2n\phi''), \quad \text{and} \\ \mathcal{A} \varphi &= \frac{2n+1}{R^2} \phi' + \frac{4\xi^2 R^2 - 2n}{R^4} \phi'' - \frac{4\xi^2}{R^4} \phi'''. \end{split}$$

In particular, we have

$$|\Delta \varphi|_{t=0} \leqslant \frac{C_0}{R^2},\tag{8}$$

where  $C_0 = 8 \|\phi''\|_{L^{\infty}} + 2n \|\phi'\|_{L^{\infty}}$ . Using the fact that  $\varphi$  satisfies Lemma 2, we obtain, for all  $R \gg 1$ , that

$$\begin{split} |\mathcal{A}\varphi|^{q} &= \left|\frac{2n+1}{R^{2}}\phi' + \frac{4\xi^{2}R^{2} - 2n}{R^{4}}\phi'' - \frac{4\xi^{2}}{R^{4}}\phi'''\right|^{q},\\ &\leqslant \frac{C_{n,q}}{R^{2q}}(|\phi'|^{q} + |\phi''|^{q} + |\phi'''|^{q}),\\ &\leqslant \frac{C_{1}}{R^{2q}}\varphi^{q-1}, \end{split}$$

where  $C_1 = C_{n,q}C_{\phi}$ .

Step 3. We perform the polar integration  $dxdt = R^{n+2}\xi^{n-1}\mathrm{d}\xi d\omega\mathrm{d}\tau$  to get that

$$\begin{split} &\iint_{K'} \frac{|\mathcal{A}\varphi|^q}{(|x|^{\sigma} \varphi)^{q-1}} \, \mathrm{d}x \, \mathrm{d}t \\ &= \iiint_{K'} \frac{C_1}{(R\xi)^{\sigma(q-1)} R^{2q}} R^{n+2} \xi^{n-1} \, \mathrm{d}\xi \, \mathrm{d}\omega \, \mathrm{d}\tau \\ &\leqslant \frac{C_1 \omega_n}{R^e} \iint_{1\leqslant \tau+\xi^2\leqslant 2} \xi^{\alpha-1} \, \mathrm{d}\xi \, \mathrm{d}\tau, \ (\omega_n = \left|S^{n-1}\right|) \\ &\leqslant \frac{2C_1 \omega_n}{R^e} \int_0^{\sqrt{2}} \xi^{\alpha-1} \, \mathrm{d}\xi, \end{split}$$

where  $\alpha = (n/(p-1))(p-1-\sigma/n)$  and  $e = (n/(p-1))(1+(\sigma+2)/n-p)$ . Since  $p > 1+((\sigma^+)/n) \ge 1+\sigma/n$ , we have  $\alpha > 0$ . Hence

$$\iint_{K'} \frac{|\mathcal{A}\varphi|^q}{(|x|^{\sigma} \varphi)^{q-1}} \leqslant \frac{M}{R^e},\tag{9}$$

where  $M = 2^{1+\alpha/2}C_1\omega_n$ . Plugging (8) and (9) in (7) yields

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$$\iint_{K} V u^{p} \varphi 
\leq \frac{M c_{0}^{1-q}}{R^{e}} + q \int_{\mathbb{R}^{n}} u_{0} (\Delta \varphi - \varphi)|_{t=0} \quad (10) 
\leq \frac{M c_{0}^{1-q}}{R^{e}} + q \int_{\mathbb{R}^{n}} u_{0} |\Delta \varphi|_{t=0} 
\leq \frac{M c_{0}^{1-q}}{R^{e}} + \frac{q C_{0}}{R^{2}} ||u_{0}||_{L^{1}}. \quad (11)$$

**Step 4.** Now consider the case  $p < 1 + (\sigma + 2)/n$ . It is obvious that e > 0. Since  $\varphi \equiv 1$  on  $\{(x, t) : 0 \leq t + |x|^2 \leq R^2\} \subset K$ , it follows that

$$\iint_{0 \leqslant t+|x|^2 \leqslant R^2} V u^p \leqslant \frac{M c_0^{1-q}}{R^e} + \frac{q C_0}{R^2} \|u_0\|_{L^1},$$

for all  $R \gg 1$ . As e > 0, the right-hand side converges to 0 as  $R \to \infty$ . Hence  $\iint_{Q_{\infty}} Vu^p = 0$  which implies  $u \equiv 0$  contradicting the non-triviality of u.

**Step 5.** For the case  $p = 1 + (\sigma + 2)/n$ , the righthand side of (11) is bounded as  $R \to \infty$ . Hence

$$\iint_{Q_{\infty}} V u^p \, \mathrm{d}x \, \mathrm{d}t < \infty.$$

Therefore  $Vu^p$  is integrable. By (5), (8), (9), and Hölder's inequality, we have

$$\begin{aligned}
\iint_{K} V u^{p} \varphi \\
\leqslant \iint_{K'} u \left| \mathcal{A} \varphi \right| + \frac{qC_{0}}{R^{2}} \left\| u_{0} \right\|_{L^{1}} \\
\leqslant \left( \iint_{K'} V u^{p} \varphi \right)^{1/p} \left( \iint_{K'} \frac{\left| \mathcal{A} \varphi \right|^{q}}{(V\varphi)^{q-1}} \right)^{1/q} \\
&+ \frac{qC_{0}}{R^{2}} \left\| u_{0} \right\|_{L^{1}} \\
\leqslant \frac{M^{1/q}}{c_{0}^{q-1}} \left( \iint_{K'} V u^{p} \varphi \right)^{1/p} + \frac{qC_{0}}{R^{2}} \left\| u_{0} \right\|_{L^{1}} \\
\leqslant \frac{M^{1/q}}{c_{0}^{q-1}} \left( \iint_{K'} V u^{p} \right)^{1/p} + \frac{qC_{0}}{R^{2}} \left\| u_{0} \right\|_{L^{1}}. (12)
\end{aligned}$$

Since  $K' \subset \{(x,t) : R^2 \leq t^2 + |x|^2 \leq 2R^2\}$ , the integrability of  $Vu^p$  implies  $\iint_{K'} Vu^p \to 0$  as  $R \to \infty$ . Therefore, by letting  $R \to \infty$ , we obtain from (12) that

$$\iint_{Q_{\infty}} V u^p = 0,$$

which implies  $u \equiv 0$ , and again a contradiction.  $\Box$ 

For the next result, we will show that when  $p > 1 + (\sigma + 2)/n$ , weak solutions to the Cauchy problem (1) blow up in a finite time if  $u_0$  is large enough.

**Theorem 2** Let  $p > 1 + (\sigma + 2)/n$ . If  $u_0 \in L^1(\mathbb{R}^n)$  is sufficiently large in the sense that there exists  $R_0 \ge C_0^{1/2}$ , where  $C_0$  is given in (8) depending on  $\phi$  from Lemma 2, such that

$$\int_{B_{R_0}(0)} u_0(x) \ge \max\left\{\frac{3}{4} \|u_0\|_{L^1}, \frac{4Mc_0^{1-q}}{qR_0^e}\right\},\$$

then every weak solution u to the Cauchy problem (1) blows up in a finite time.

*Proof*: Again, we will prove by contradiction and therefore assume that the global weak solution u exists. Set  $R = R_0 \gg 1$  in the proof of the preceding theorem. Since  $R_0 \ge C_0^{1/2}$ , (8) can be reduced to

$$|\Delta\varphi(x,0)| \leqslant 1.$$

In (10), which is true for all cases of  $p > 1 + (\sigma^+)/n$ , the second term on the right-hand side can be estimated by

$$\begin{split} \int_{\mathbb{R}^n} u_0(x) [\Delta \varphi(x,0) - \varphi(x,0)] \, \mathrm{d}x \\ &\leqslant \int_{|x| \ge R_0} u_0(x) - \int_{|x| \le R_0} u_0(x) \\ &= \|u_0\|_{L^1} - 2 \int_{|x| \le R_0} u_0(x) \\ &\leqslant -\frac{1}{2} \|u_0\|_{L^1} \leqslant -\frac{1}{2} \int_{|x| \le R_0} u_0(x). \end{split}$$

Hence

$$\iint_{K} V(x) u^{p} \varphi \leq M c_{0}^{1-q} R^{-e} - \frac{q}{2} \int_{|x| \leq R_{0}} u_{0}(x)$$
$$\leq M c_{0}^{1-q} R^{-e} - 2M c_{0}^{1-q} R^{-e}$$
$$= -M c_{0}^{1-p'} R^{-e} < 0,$$

which is absurd because  $V, u \ge 0$ . Therefore there is no global weak solution to (1).

**Remark 2** Examples of  $u_0$  satisfying the conditions of Theorem 2 are

$$u_0 = a \chi_{B_{R_0}(0)}$$

where  $R_0 \ge C_0^{1/2}$  and a is a constant satisfying

$$a \geqslant \frac{4Mc_0^{1-q}}{\omega_n q} C_0^{-(n+e)/2}$$

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